

Particle traps and β -correlations to probe weak processes in nuclei

- 1. Introduction**
- 2. Brief formalism**
- 3. Status and plans**

Weak Interactions in Nuclei and Astrophysics: Standard Model and Beyond
ECT* Trento, 16 - 21 June 2003
N. Severijns
Kath. Univ. Leuven

1. Introduction – The Standard Model and beyond

Standard Model :

- works well, but still many problems as well
- SM ‘low’ energy (~ 200 GeV) approximation of more fundamental theory
- search for physics beyond SM (e.g. neutrino oscillations !!)

Testing the Standard Model in the sector of the weak interaction :

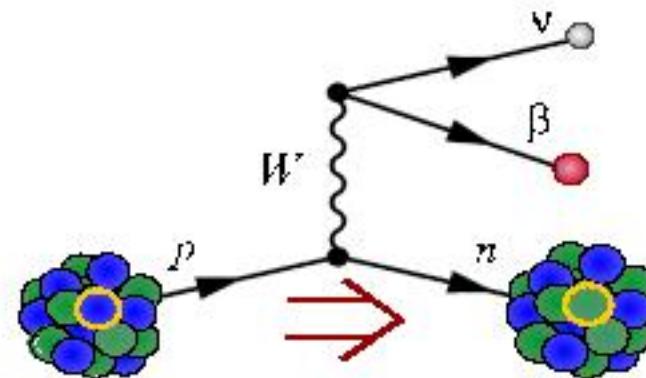
- at high energy colliders (CERN, Fermilab, DESY, ...)
- in neutrino physics (SuperKamiokande, AMANDA, ...)
- in atomic physics (e.g. parity violation)
- in nuclear beta decay (correlations, ft-values)
- ...

2. STRUCTURE OF THE WEAK INTERACTION IN NUCLEAR BETA DECAY

Fermi, 1933 : 4-fermion interaction

Lee & Yang, 1956; Wu et al., 1957 :
parity violation

→ general Lorentz invariant
4-fermion interaction



$$H_\beta \propto \sum_i \bar{e} O_i (C_i + C'_i \gamma_5) \nu_e \langle p | O_i | n \rangle + \text{h.c.}$$

with **i** = **S**calar, **V**ector, **T**ensor, **A**xial vector, **P**seudoscalar
and **coupling constants C_i** defining properties of the interaction types

Transition probability :

**distribution in energy, emission angle and polarization of β -particles
for allowed β -decay of polarized nuclei**

$$dW = dW_0 \xi \left\{ 1 + \frac{\bar{p} \cdot \bar{q}}{E_e E_\nu} a + \frac{\Gamma m}{E_e} b \right.$$
$$+ \bar{J} \cdot \left[\frac{\bar{p}}{E_e} A + \frac{\bar{q}}{E_\nu} B + \frac{\bar{p} \times \bar{q}}{E_e E_\nu} D \right]$$

$$+ \bar{\sigma} \cdot \left[\frac{\bar{p}}{E_e} G + \hat{p} (\bar{J} \cdot \hat{p}) Q' + \bar{J} \times \frac{\bar{p}}{E_e} R \right]$$

$$dW_0 \propto G_F^2 F(\pm Z, E_e) (E_e - E_0)^2 p E_e dE d\Omega_e d\Omega_\nu$$

with :

p, m, E_e beta-particles momentum, mass and total energy

q, E_ν neutrino momentum and total energy

J nuclear spin polarization vector

σ beta-particles spin vector

$\Gamma = \text{sqrt} [1 - (\alpha Z)^2]$, Coulomb factor

Note: contribution from the part of the Hamiltonian involving the pseudoscalar quark current $\bar{u}\gamma_5 d$ can be neglected as it vanishes in the nonrelativistic approximation of the nucleons;

also, from $BR(\pi \rightarrow e\bar{\nu}_e)/BR(\pi \rightarrow \mu\bar{\nu}_\mu)$: $C_P < 1.25 \times 10^{-4} C_A$

e.g. $\beta\nu$ -correlation coefficient :

$$\alpha\xi = M_F^2 \left[|C_V|^2 + |C'_V|^2 - |C_S|^2 - |C'_S|^2 \right] - \frac{1}{3} M_{GT}^2 \left[|C_A|^2 + |C'_A|^2 - |C_T|^2 - |C'_T|^2 \right]$$

with

$$\xi = M_F^2 \left[|C_V|^2 + |C'_V|^2 + |C_S|^2 + |C'_S|^2 \right] + M_{GT}^2 \left[|C_A|^2 + |C'_A|^2 + |C_T|^2 + |C'_T|^2 \right]$$

→ independent of nuclear matrix elements if pure F or GT transition is used

(radiative corrections : order 10^{-4} to 10^{-5} typically,

recoil corrections : order 10^{-3} to 10^{-4} ;

both to be addressed in more detail if order 10^{-3} precision is reached)

$\beta\nu$ -correlation

$$W(\theta) = \frac{dW_0 \xi \left\{ 1 + \frac{\bar{p} \cdot \bar{q}}{E_e E_\nu} a + \frac{\Gamma m}{E_e} b \right\}}{dW_0 \xi \left\{ 1 + \frac{\Gamma m}{E_e} b \right\}} \Rightarrow W(\theta) = 1 + \frac{\bar{p} \cdot \bar{q}}{E_e E_\nu} \tilde{a}$$

with $\tilde{a} \equiv \frac{a}{1 + \frac{\Gamma m}{E_e} b}$

$$a_F \equiv 1 - \frac{|C_S|^2 + |C_S'|^2}{|C_V|^2}$$

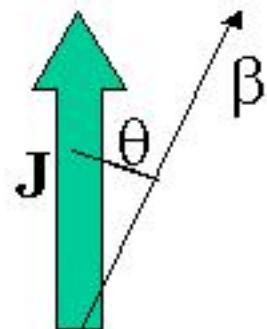
$$a_{GT} \equiv -\frac{1}{3} \left[1 - \frac{|C_T|^2 + |C_T'|^2}{|C_A|^2} \right]$$

$$b_F \equiv R e \frac{C_S + C_S'}{C_V}$$

$$b_{GT} \equiv R e \frac{C_T + C_T'}{C_A}$$

(assuming maximal P-violation and T-invariance for V- and A-interactions)

β -asymmetry



$$W(\theta) = 1 + \bar{J} \cdot \frac{\bar{p}}{E_e} \tilde{A}$$

$$\text{with } \tilde{A} \equiv \frac{A}{1 + \frac{\Gamma m}{E_e} b}$$

for a pure Gamow-Teller transition :

$$\tilde{A}_{GT} \cong \mp 1 + \frac{\alpha Zm}{p_e} \operatorname{Im} \left(\frac{C_T + C'_T}{C_A} \right) \pm \frac{\Gamma m}{E_e} \operatorname{Re} \left(\frac{C_T + C'_T}{C_A} \right)$$

(assuming maximal P-violation and T-invariance for V- and A-interactions)

In the Standard Model :

- * $C_V = 1$
- * $C_A = -1.26$
- * $C_V' = C_V$ and $C_A' = C_A$

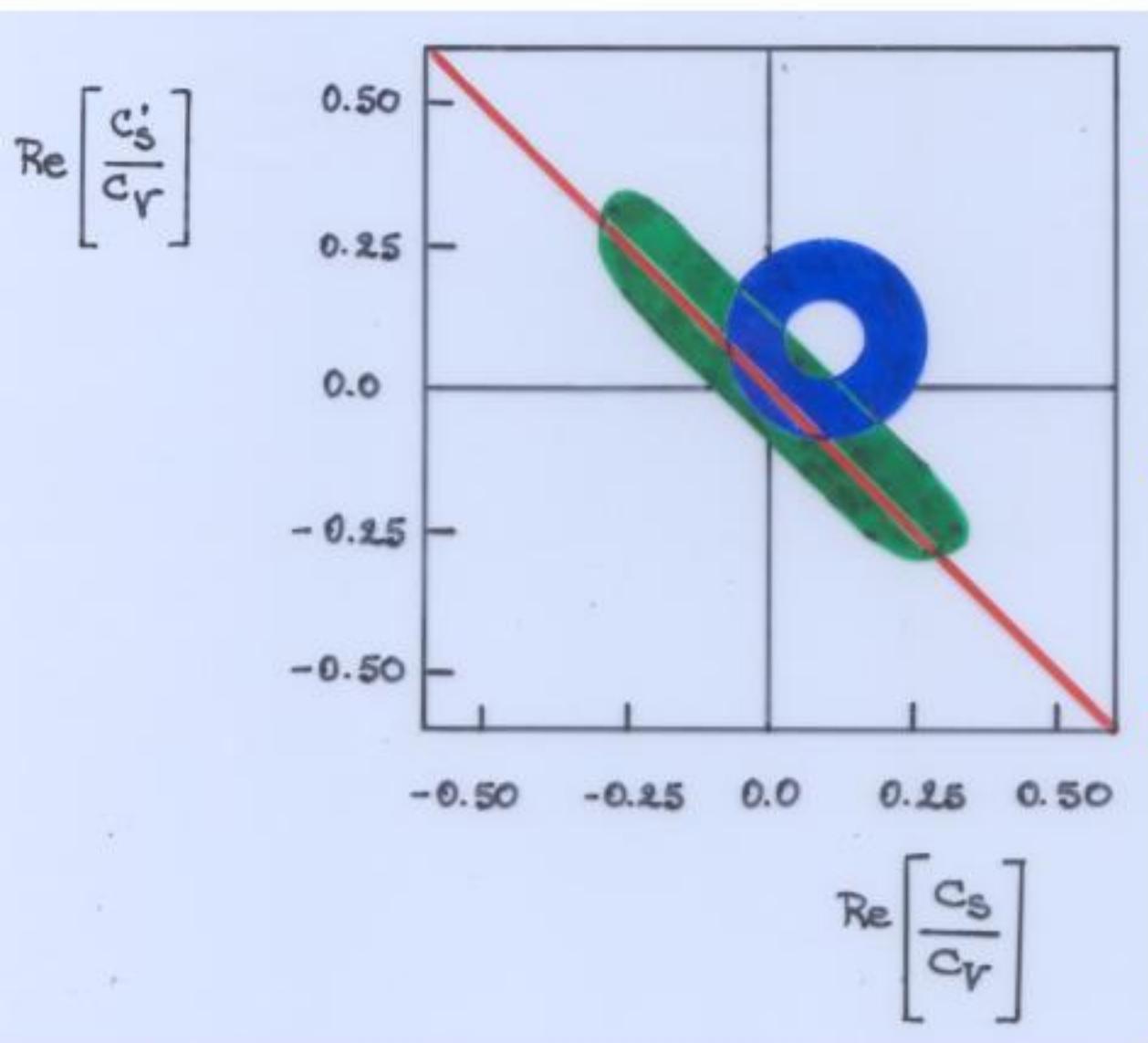
* $C_S = C_S' = C_T = C_T' = C_P = C_P' \equiv 0$

experimental upper limits:

$$\left| \frac{C_T}{C_V} \right| < 0.13 \quad \left| \frac{C_S}{C_V} \right| < 0.08 \quad 95\% \text{CL}$$

- * **no time reversal violation** (except for the CP-violation described by the phase in the CKM quark-mixing matrix)

Present β -decay limits on scalar currents



green:

neutron decay experiments,
positron polarization (unpol. nuclei)
Fierz interference term in ^{22}Na & ^6He

blue :

$\beta\nu$ -correlation, ^{32}Ar

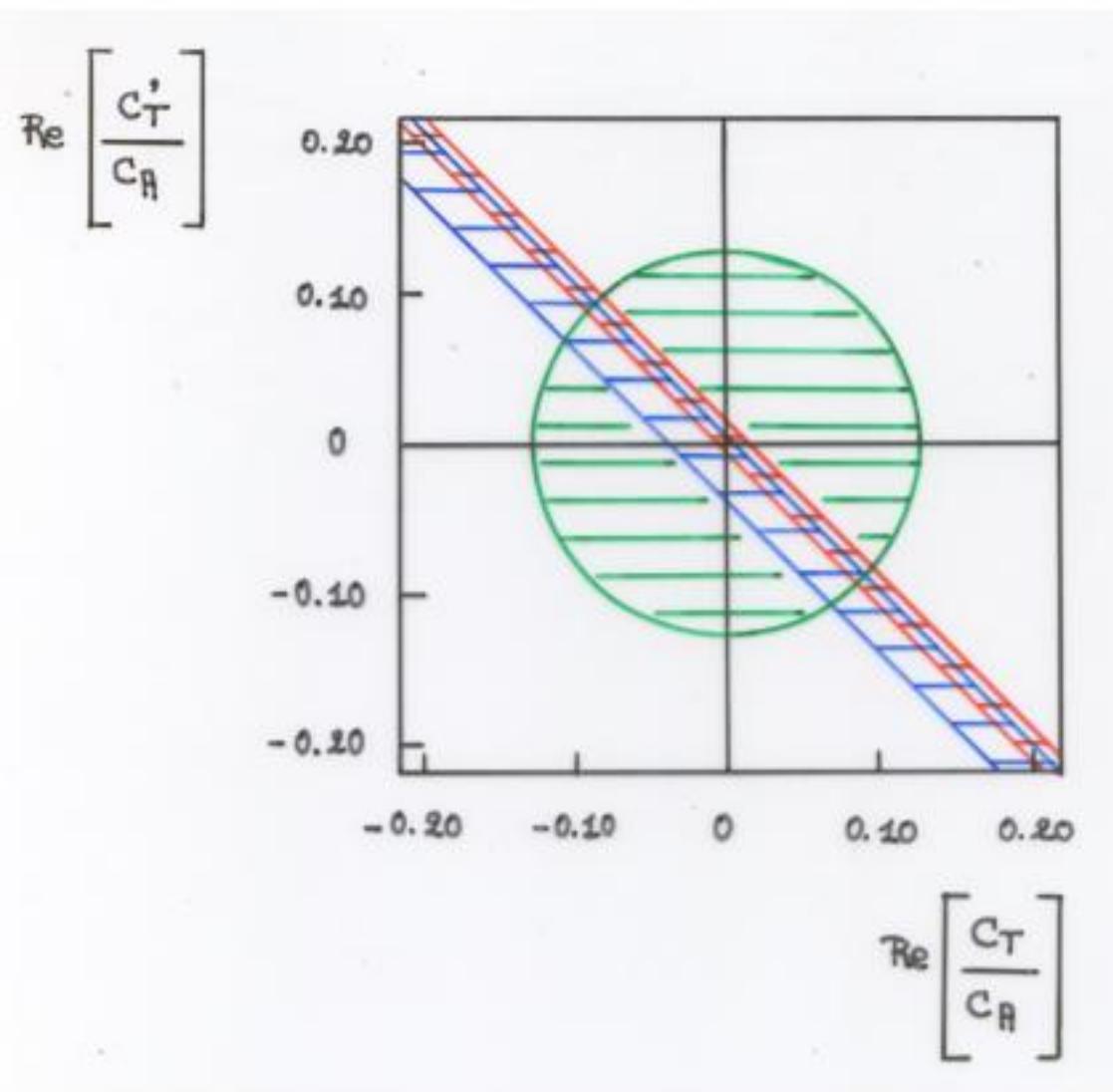
Adelberger et al., Phys. Rev. Lett. 83 (1999)

red :

Ft-value of superallowed Fermi transitions

L.S. Towner, J.C. Hardy, J.Phys. G: Nucl. Part. Phys. 29 (2003) 197

Present β -decay limits on tensor currents



green:

$\beta\nu$ -correlation, ${}^6\text{He}$

Johnson et al., Phys.Rev. 132 (1963) 1149

blue :

positron polarization, ${}^{107}\text{In}$

Severijns et al., Nucl. Phys. A629 (1998) 423c

red :

all other relevant experiments

P.A. Quin et al., Phys. Rev. D47 (1993) 1247

Limits (95 % CL) on possible new bosons for S- and T-interactions, from high-energies:

mass limit for H^\pm (charged Higgs) : $> 71.5 \text{ GeV}$ (LEP)

**mass limit for leptoquarks : $> 242 \text{ GeV}$ (from pair production; combined CDF-D0)
: $> 290 \text{ GeV}$ (from single production; H1)**

Sensitivity of beta-neutrino correlation (95 % CL) :

$\Delta a = 0.01 \rightarrow$ sensitive to masses of new bosons of $\sim (0.01)^{-1/4} M_W \approx 215 \text{ GeV/c}^2$

$\Delta a = 0.005 \rightarrow$ sensitive to masses of new bosons of $\sim (0.005)^{-1/4} M_W \approx 255 \text{ GeV/c}^2$

$\Delta a = 0.002 \rightarrow$ sensitive to masses of new bosons of $\sim (0.002)^{-1/4} M_W \approx 320 \text{ GeV/c}^2$
("handwaving")

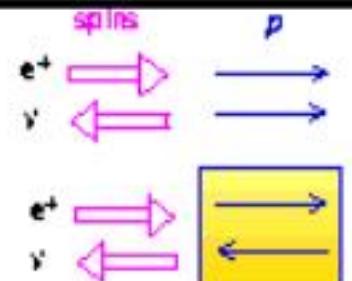
Also, it always makes sense to carry out tests of the SM in different energy domains.

(e, v) correlation in the β decay of ^{32}Ar

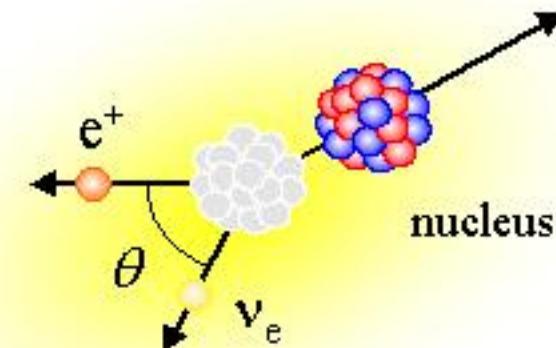
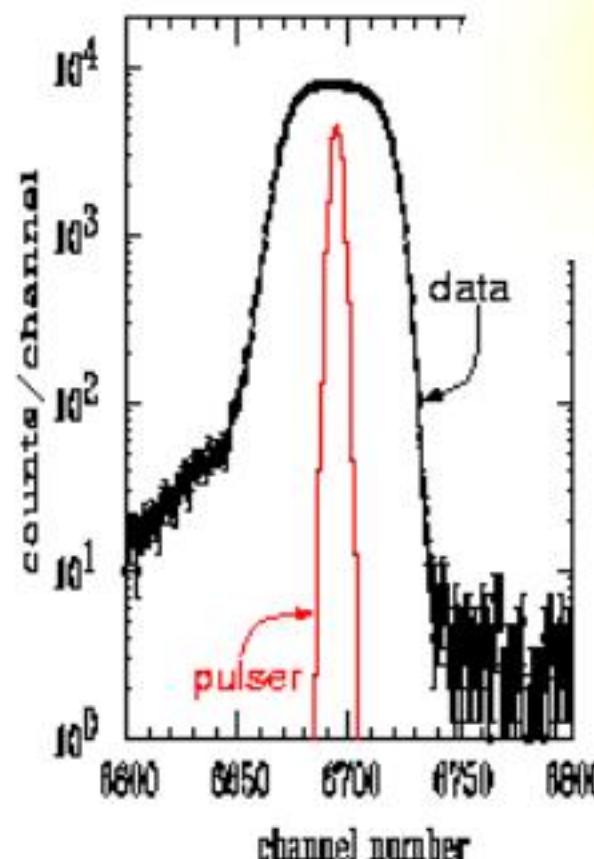
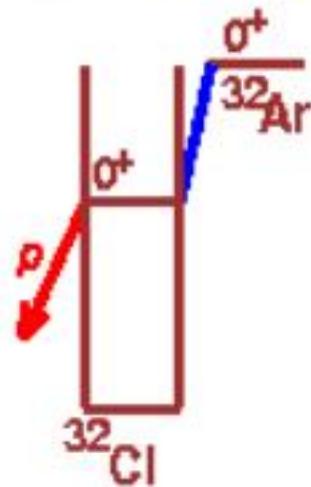
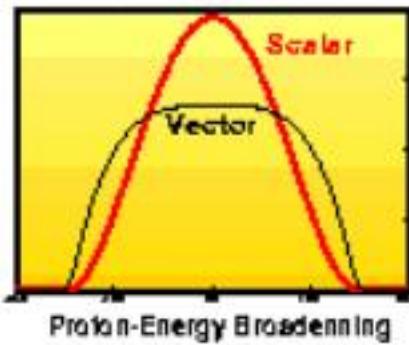
E.G. Adelberger, C. Ortiz, A. Garcia, H.E. Swanson, M. Beck, O. Tengblad,
 M.J.G. Borge, I Martel, H. Bichsel, and the ISOLDE collaboration,
 PRL 83, 1299 (1999).

$$\frac{dW/d\Omega(\text{Vector})}{dW/d\Omega(\text{Scalar})} = 1 + \frac{\vec{P}_e \cdot \vec{P}_v}{E_e E_v}$$

$$= 1 - \frac{\vec{P}_e \cdot \vec{P}_v}{E_e E_v}$$



We measure the (e, v) correlation by detecting the Doppler broadening of the β -delayed proton:



$$\tilde{a} = 0.9989 \pm 0.0052(\text{stat}) \pm 0.0039 \text{ (syst)} \rightarrow \tilde{a} = 1.0050 \pm 0.0052(\text{stat}) \pm (\text{syst})$$

after measurement of ^{32}Ar mass (K. Blaum et al., subm. to PRL)

3. Traps for weak interaction physics

Ion and atom traps
provide ideal sources for
weak interaction tests
in nuclear beta decay :

- sample is isotopically **pure**
- localized in **small volume**
- source **scattering negligible**
- atoms/ions **decay at rest**
- potential for polarized sample

Traps for weak interaction physics :

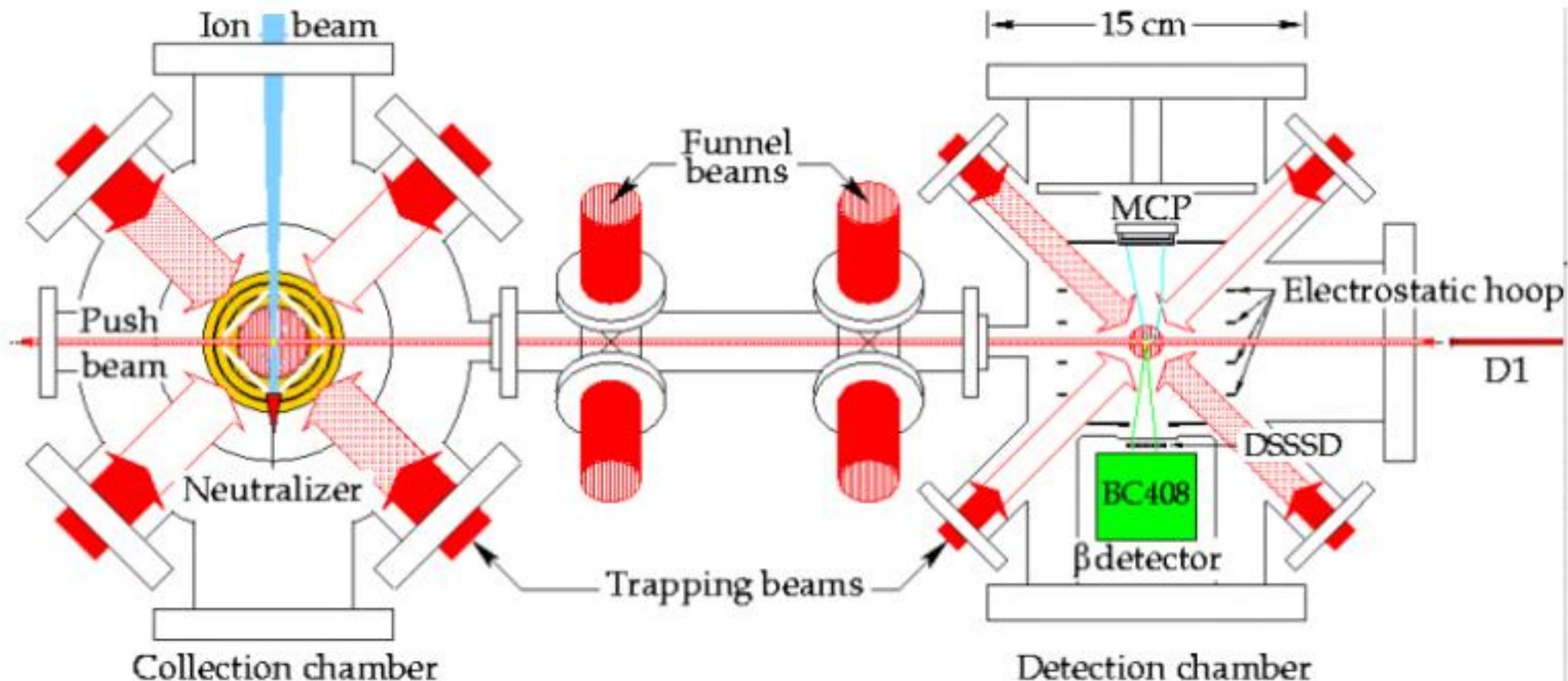
1. Atom traps :

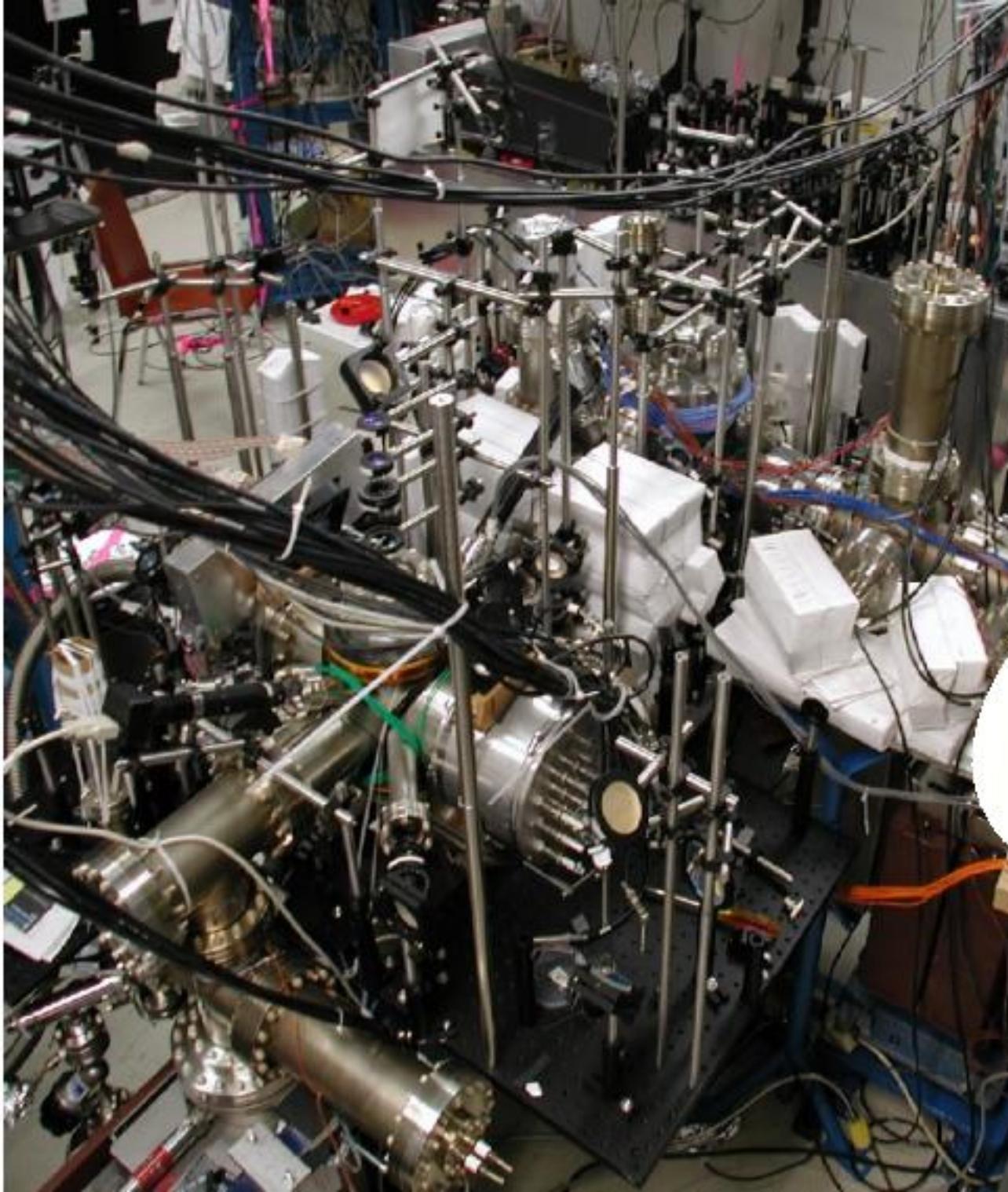
- TRIUMF-ISAC, ^{38m}K , $\beta\nu$ -correlation (J. Behr et al.)
A. Gorelov et al., Hyperfine Interactions 127 (2000) 373
- LBNL & UC Berkeley, ^{21}Na , $\beta\nu$ -correlation (S.J. Freedman et al.)
N. Scielzo, Ph. D. Thesis (2003)
- LANL Los Alamos, ^{82}Rb , β -asymmetry (D. Vieira et al.)
S.G. Crane et al., Phys. Rev. Lett. 86 (2001) 2967
- KVI-Groningen, Na, Ne, Mg, D-coefficient (K. Jungmann et al.)
- ...

2. Ion traps :

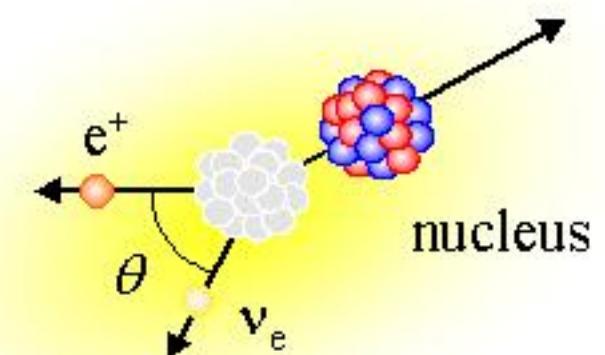
- LPC-Caen, ^6He , $\beta\nu$ -correlation (O. Naviliat-Cuncic et al.)
P. Delahaye et al., Hyperfine Interactions 132 (2001) 479
- Leuven, ^{35}Ar , $\beta\nu$ -correlation (N. Severijns et al.)
D. Beck et al., Nucl. Inst. Methods Phys. Res., A 503 (2003) 567
- ISOLTRAP-CERN, mass for $0+ \rightarrow 0+$ decays (H.-J. Kluge et al.)
- ...

TRIUMF Neutral Atom Trap

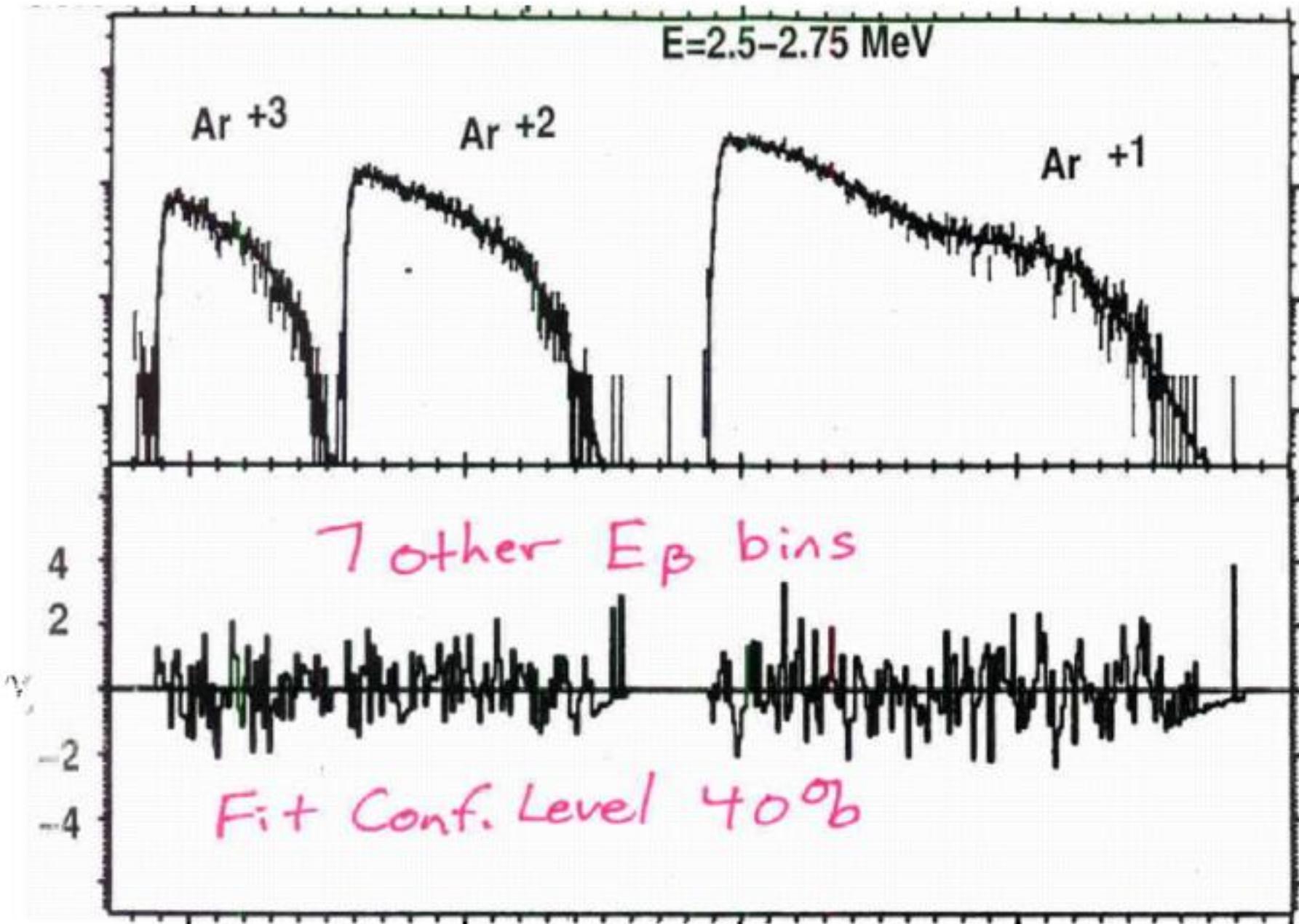




TRIUMF
neutral atom trap



J. Behr et al.



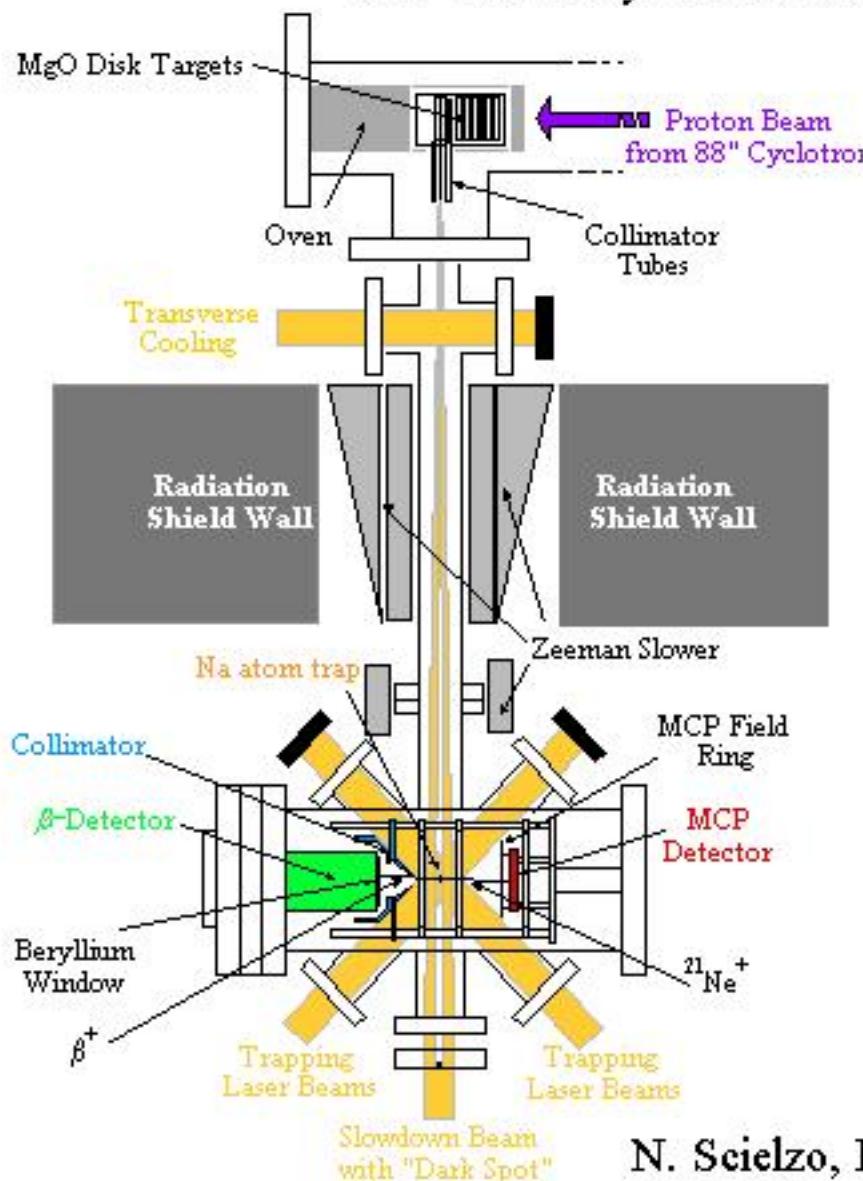
First result available in autumn 2003 (expected)



Measuring the $\beta-\nu$ Angular Correlation in Magneto-Optically Trapped ^{21}Na

88-inch
Cyclotron

S. J. Freedman, B. K. Fujikawa, N. D. Scielzo, and P. A. Vetter
UC Berkeley and Nuclear Science Division, LBNL

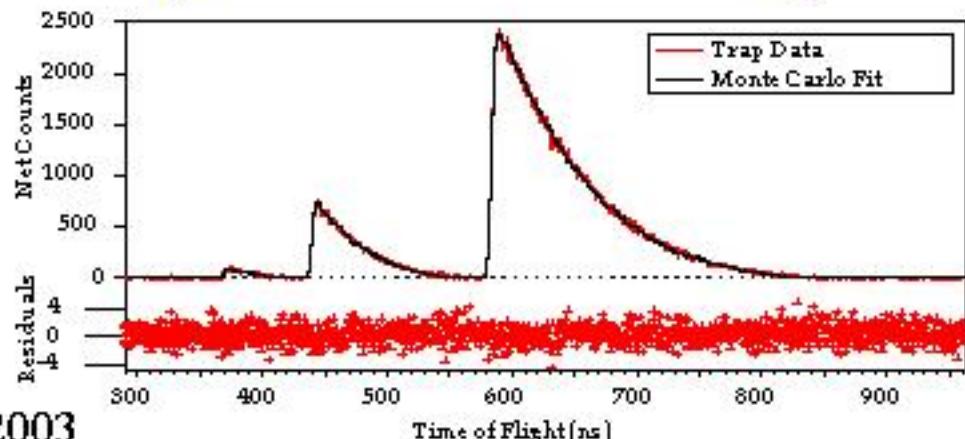
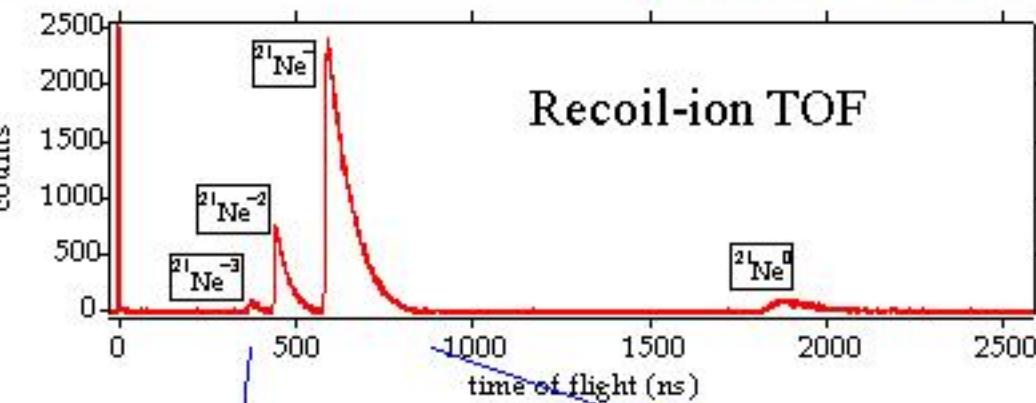


Result:

$$\alpha_{exp} = 0.5243 \pm 0.0092$$

$$\alpha_{SM} = 0.558 \pm 0.003$$

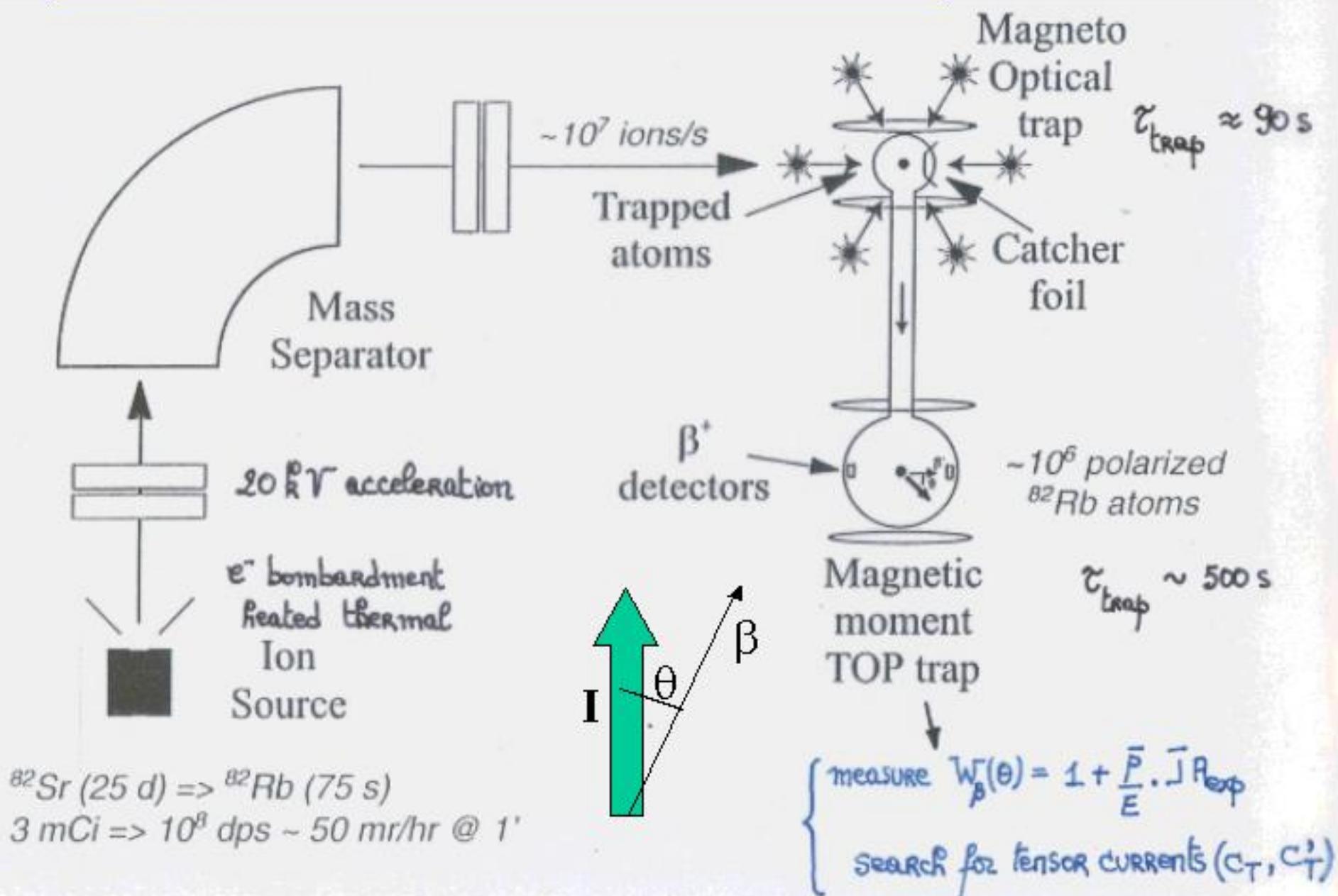
500,000
trapped atoms



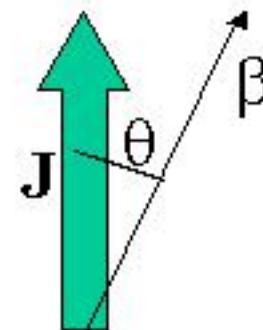
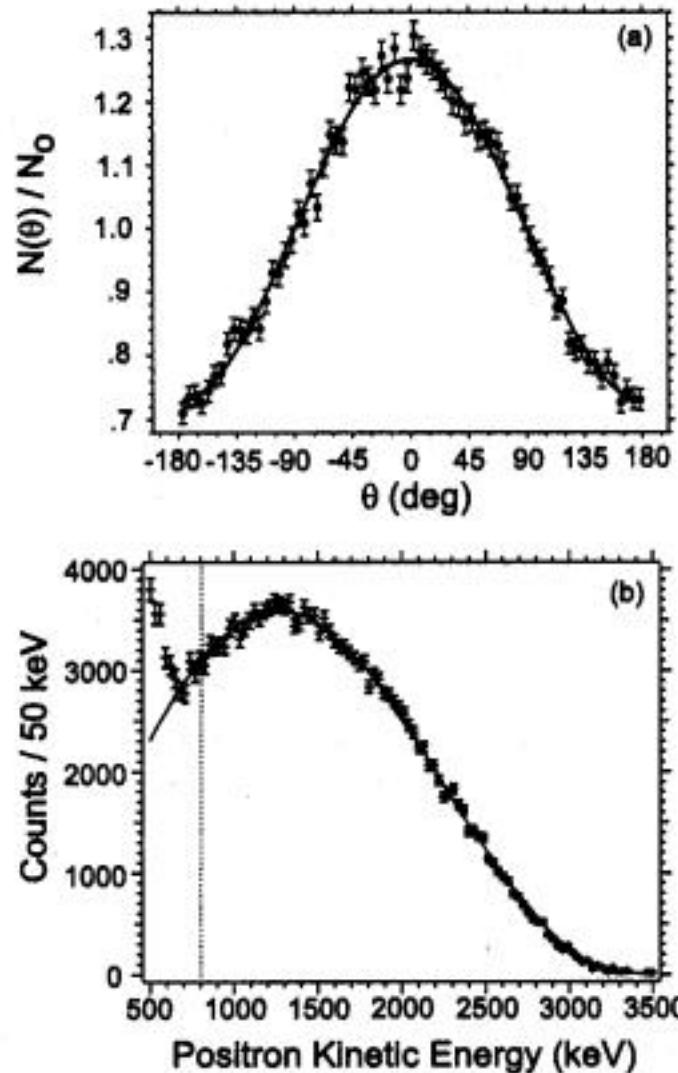
N. Scielzo, Ph.D. May 2003

The atom trap at Los Alamos - LANL

D. Vieira et al.



angular distribution ↓



D.J. Vieira et al., Hyp.Int. 127 (2000) 387

S.G. Crane et al., Phys. Rev. Lett. 86 (2001) 2967 (first results)

positron energy spectrum ↑

The LPC-Caen Paul trap

The RFQ cooler and decelerating electrode in the high voltage platform

gas inlet

insulator

Spiral Beam

$10^{9.6} \text{He/s at } 30 \text{ keV}$

emittance: $80\pi \text{ mm.mrad}$

decelerating electrode

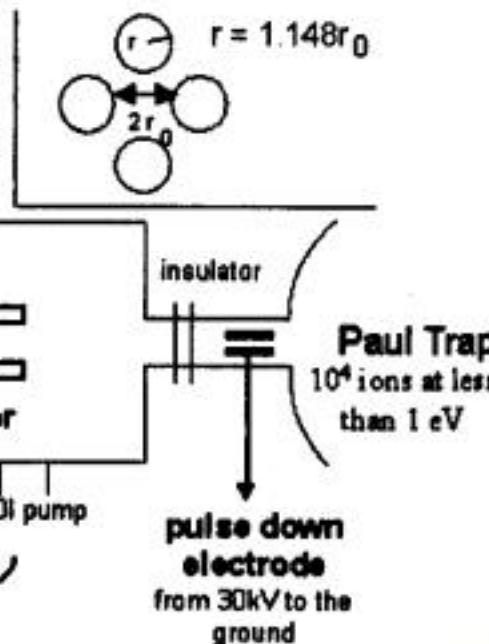
cooler

buncher

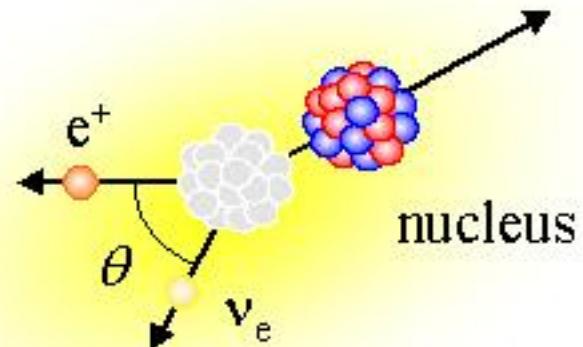
1000l pump

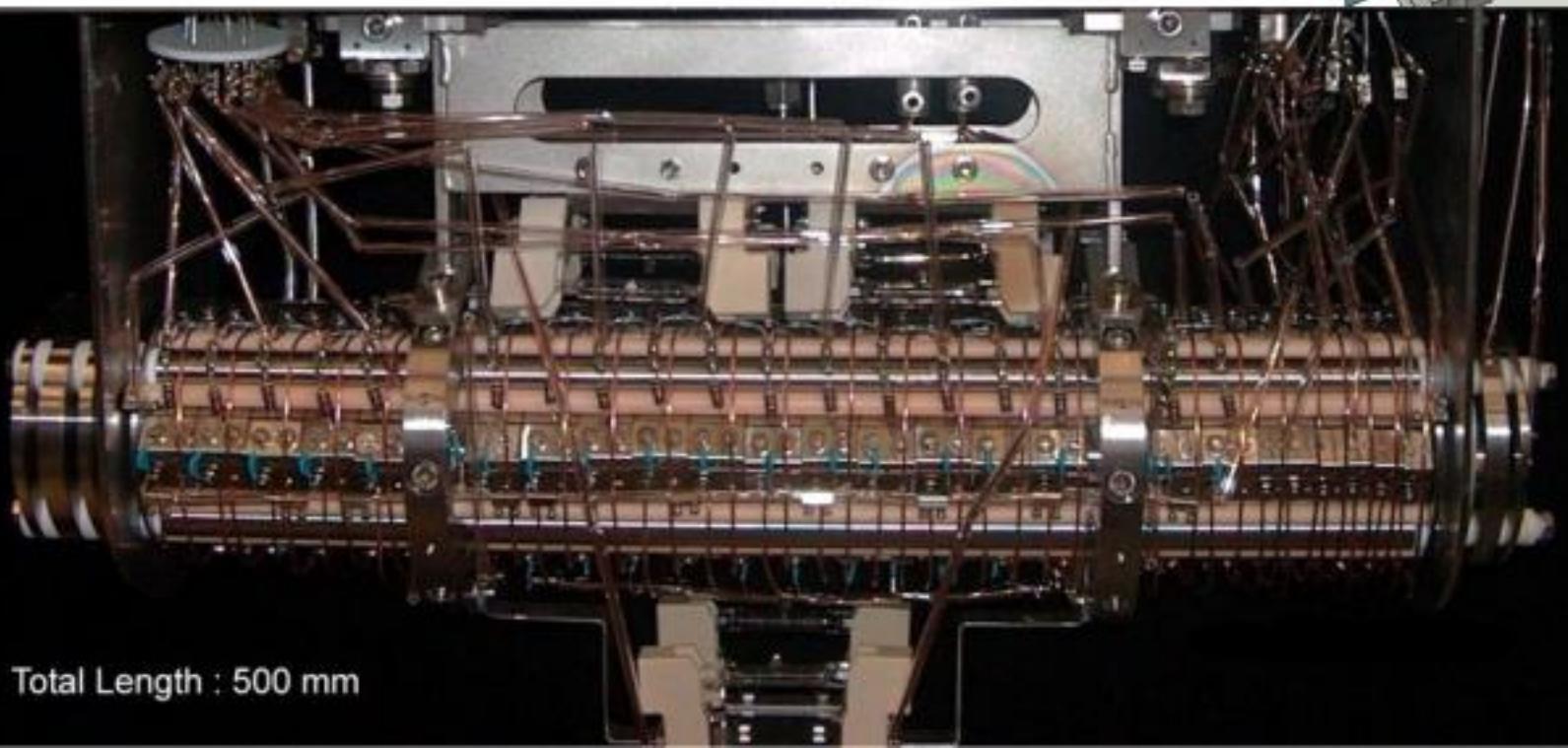
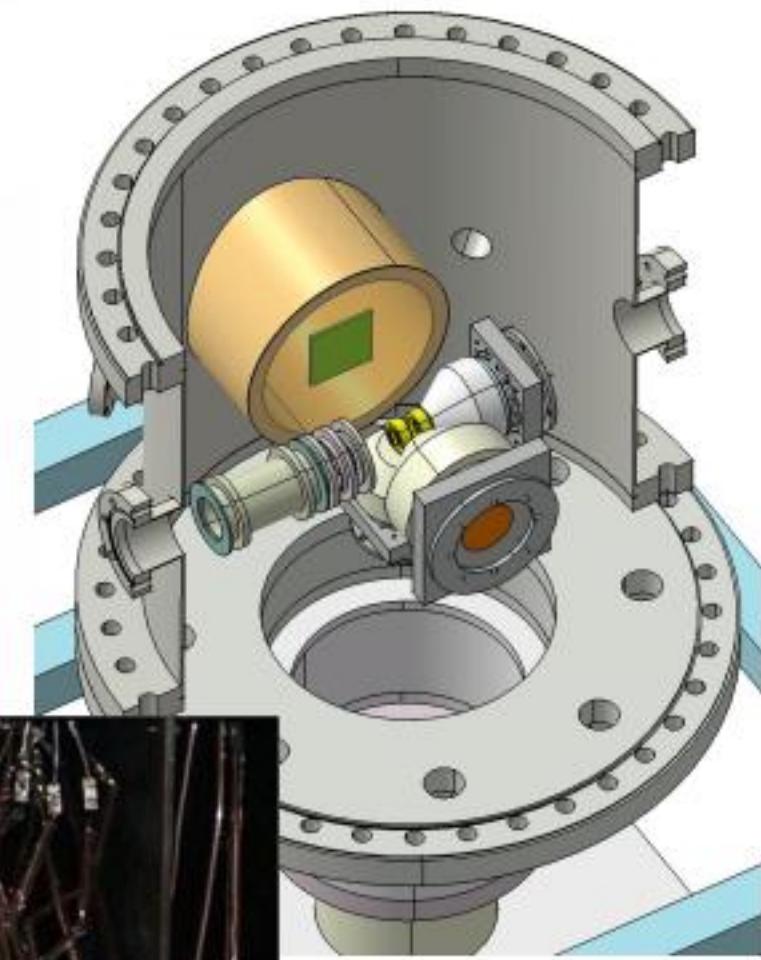
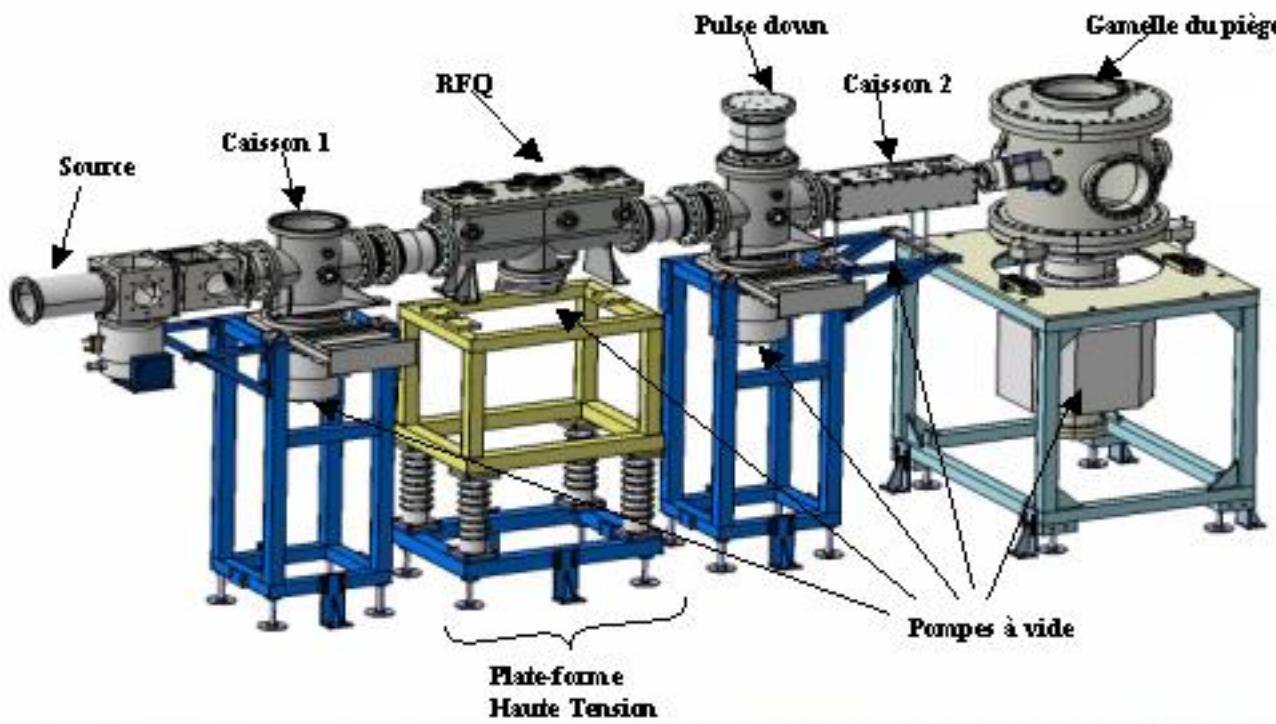
1000l pump

high voltage platform
30kV



- measure $\beta\nu$ -correlation with ${}^6\text{He}$
- cool ${}^6\text{He}$ with hydrogen gas
- to be installed at SPIRAL





Total Length : 500 mm

First experiment
expected in 2004

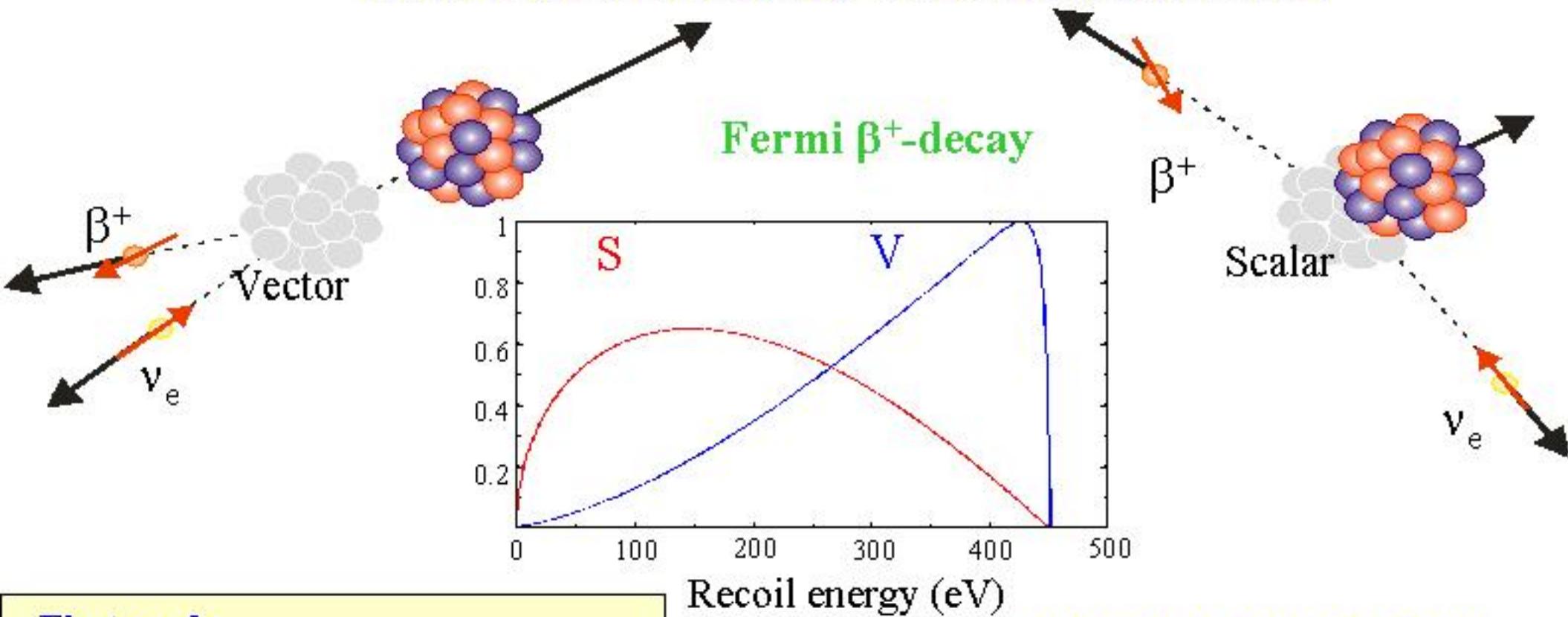
LPC Paul trap set-up
O. Naviliat-Cuncic et al.



WITCH – Weak Interaction Trap for CHarged particles

ISOLDE-CERN (K.U.Leuven, Univ. Munchen, CERN)

cooler & decay Penning trap + retardation spectrometer



First goal :

search for scalar weak interaction
by measuring
shape of recoil ion energy spectrum
after β -decay

First experiment
expected in 2004

Other physics possibilities :

- in-trap beta spectroscopy
- determination of EC/ β^+ ratios
- determination of Q_β -values
- measure charge state distributions



Detector (MCP)

Experimental set-up

Post acceleration

Energy analysis by retardation

Radial into axial motion conversion

Decay trap

Cooler trap

Vertical beam line

Horizontal beam line

Einzellens

β -detector

0.1T magnet

9T magnet

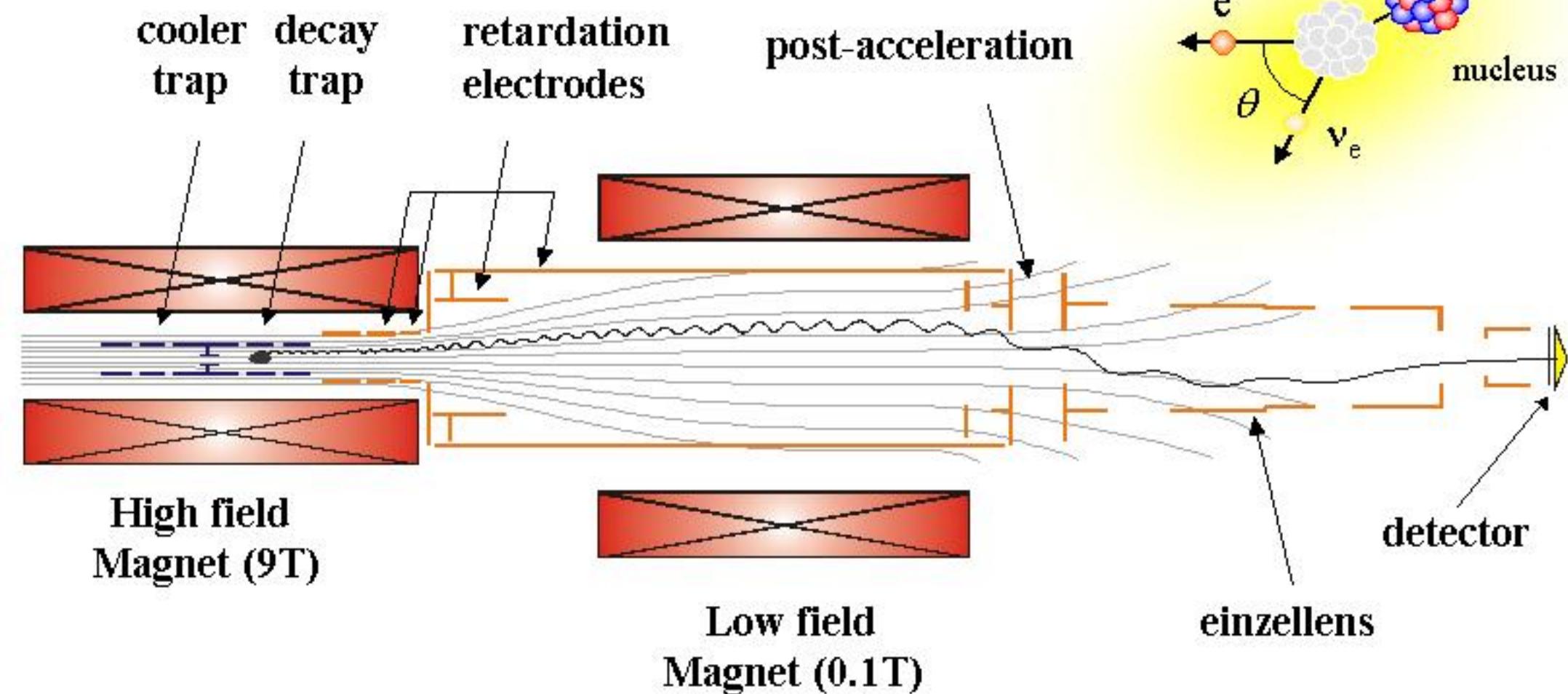
Pulsed drift tube

90° bender

REXTRAP



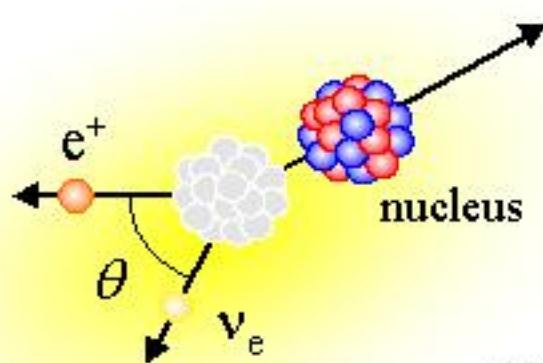
WITCH retardation spectrometer



$$\frac{E_{\perp 1}^{kin}}{E_{\perp 0}^{kin}} = \frac{B_1}{B_0} = \frac{0.1T}{9T} = 1.1\%$$



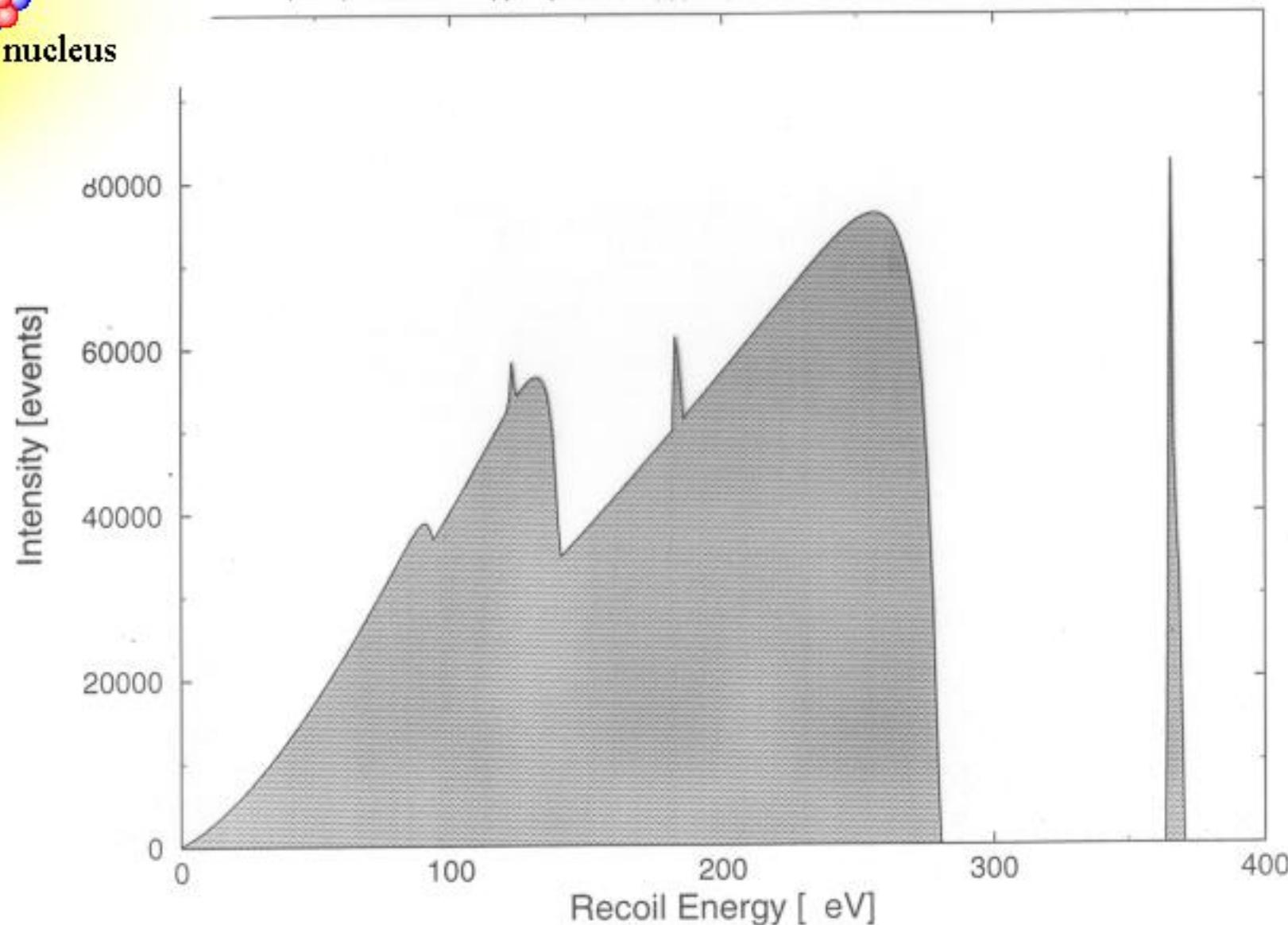
What we expect to see :

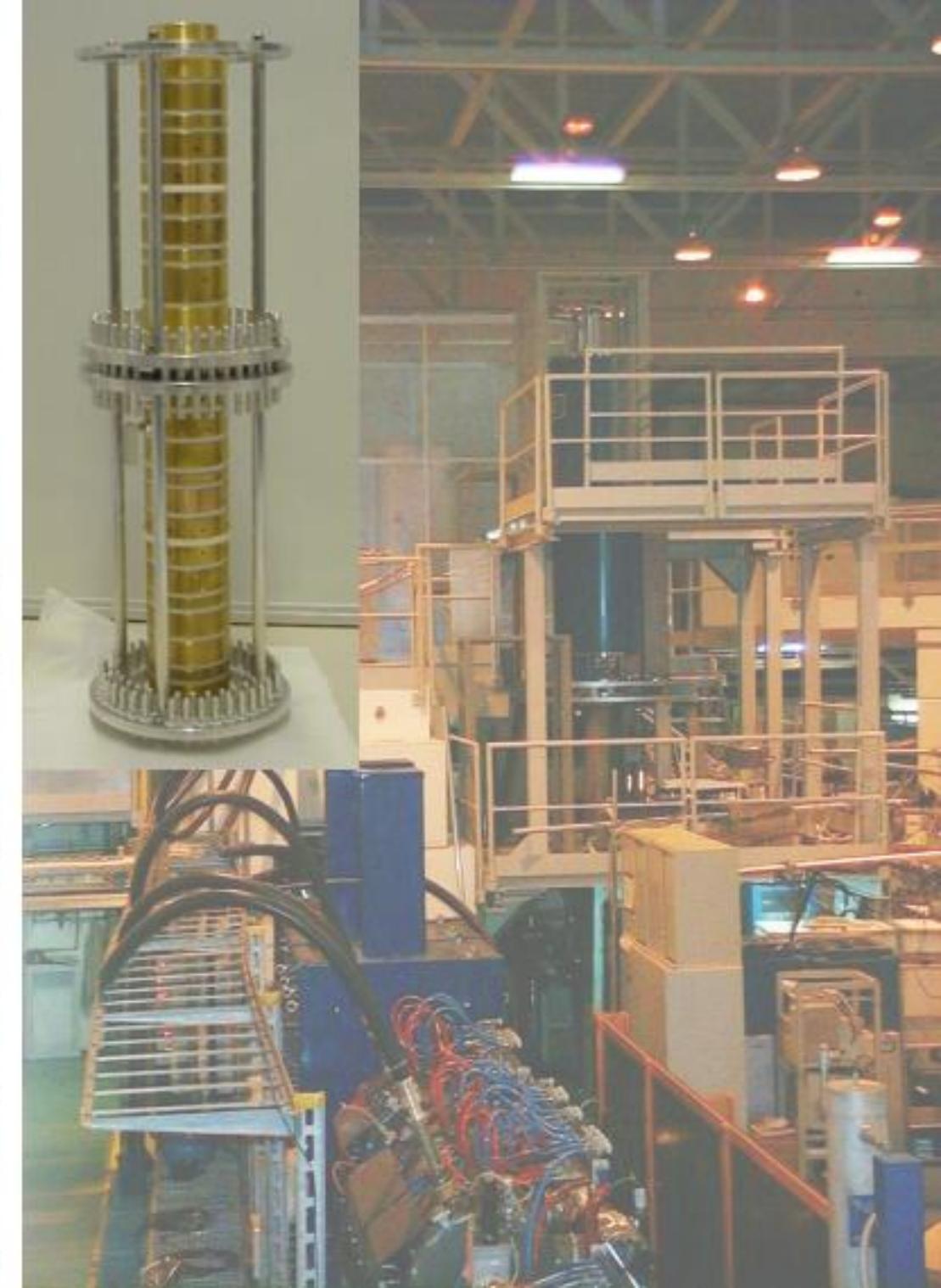


Full Recoil Energy Spectrum



$^{26}\text{Al}^m$, $a=1$, assumed $\text{BR}(q=+1)=13\%$, $\text{BR}(q=+2)=2\%$, $\text{BR}(q=+3)=0.25\%$, $\text{BR}(\text{EC}/\beta^+)=2.5\%$





Experiments not using traps (all in preparation) :

1. Neutron decay

- **aspect spectrometer** (Mainz, München, ILL) / α -coefficient
O. Zimmer et al., Nucl. Instr. Meth. A 440 (200) 548
- **PERKEO-II** (Heidelberg, ILL) / A -coefficient
H. Abele et al.

2. Nuclear β -decay

- **32Ar** (LPC-Caen, GANIL) / α -coefficient
O. Naviliat-Cuncic et al.
- **polarized nuclei** (Leuven, ISOLDE-CERN) / A -coefficient
N. Severijns et al.

3. The unitarity problem : $\sum V_{ui}^2 = V_{ud}^2 + V_{us}^2 + V_{ub}^2 = 1$?

V_{ud} :

a. *Ft* ($0^+ \rightarrow 0^+$)

$$Ft = ft(1 + \delta_R)(1 - \delta_C) = \frac{K}{2G_F^2 V_{ud}^2 (1 + \Delta_R)} \rightarrow \boxed{\sum V_{ui}^2 = 0.9969(15)}$$

b. neutron decay (τ_n and *A*)

world average values $\rightarrow \sum V_{ui}^2 = 0.9961(24)$

only most precise *A*-result $\rightarrow \sum V_{ui}^2 = 0.9917(28)$

c. pion β -decay : $\sum V_{ui}^2 = 1.0030(110)$

based on $V_{us} = 0.2196(26)$ from K_{e3} decay ($K^+ \rightarrow \pi^0 e^+ \nu_e$) (late 1970's data)

New BNL-data : $V_{us} = 0.2272(20)(7)(18)$ from K_{e3} decay (Ph.D. thesis A. Sher)

$$\rightarrow \boxed{\sum V_{ui}^2 = 1.0003(16)} \text{ (for } 0^+ \rightarrow 0^+ \text{) and } \sum V_{ui}^2 = 0.9995(25) \text{ (for neutron) !!!}$$

→ New BNL data for V_{us} (need to be confirmed) solve unitarity problem !!